

## NONLINEAR PROPAGATION OF ULTRASOUND IN LIQUID MEDIA

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It is well known that acoustic phenomena are fundamentally non-linear, though a large class of acoustical phenomena can be dealt with by linearizing the equations of motion leading to acceptable and usable solutions(1,2). This is, of course, very convenient as the nonlinear relations are extremely difficult to deal with. However, it is necessary on occasion to have more details of the propagation processes, e.g., for a more profound understanding of the phenomena, than can be obtained from the linear theories. Thus, a program has been initiated to examine this question and preliminary measurements give an indication of the degree of non-linearity of biological media.

The equation of motion or the wave equation for acoustic phenomena is obtained by invoking three constitutive equations, viz. an equation of continuity, a dynamical equation, and an equation of state of the medium in which the propagation is to take place. In Lagrangian coordinates, these are respectively:

$$\rho = \rho_0 \left(1 + \frac{\partial \xi}{\partial x}\right)^{-1}, \quad p \dot{\xi} = - \frac{\partial p}{\partial x}, \quad \text{and} \quad c^2 = \frac{\partial p}{\partial \rho} \quad (1)$$

Here,  $\rho$  and  $\rho_0$  are the disturbed and undisturbed density, respectively,  $\xi$  is the particle displacement of the medium,  $p$  is the sound pressure, and  $x$  is the Lagrangian coordinate. Combining these equations leads to:

$$\dot{\xi} = c^2 / \left(1 + \frac{\partial \xi}{\partial x}\right)^2 \xi'' ,$$

equation of motion obtained without approximation. It is seen that if a Hooke's law relationship for the equation of state is used, such that  $\partial p/\partial \rho$  is constant, and if the displacement of the wave is sufficiently small such that  $\partial \xi/\partial x \ll 1$ , the ordinary lossless wave equation is obtained, viz.  $\xi = c^2 \xi''$ .

Now consider the situation wherein the displacement amplitude is negligible and the equation of state is more complex than the Hooke's law relationship. The second point can be accomplished by considering the equation of state to be expanded in a Taylor series for the isentropic case such that:

$$\begin{aligned} \rho - \rho_0 &= \left(\frac{\partial \rho}{\partial p}\right)_{S, \rho} (\rho - \rho_0) + \frac{1}{2!} \left(\frac{\partial^2 \rho}{\partial \rho^2}\right) (\rho - \rho_0)^2 + \dots \\ &+ \frac{1}{3!} \left(\frac{\partial^3 \rho}{\partial \rho^3}\right) (\rho - \rho_0)^3 + \dots \end{aligned} \quad S, \rho = \rho_0$$

convenient to rewrite this in series form

$$\rho - \rho_0 = A s + \frac{B}{2!} s^2 + \frac{C}{3!} s^3 + \dots,$$

$$A = \rho_0 \left(\frac{\partial \rho}{\partial p}\right)_{S, \rho} = \rho_0 c_0^2;$$

$$B = \frac{2}{\rho_0} \left(\frac{\partial^2 \rho}{\partial \rho^2}\right)_{S, \rho} = \rho_0 \quad ;$$

$$C = \frac{3}{\rho_0} \left(\frac{\partial^3 \rho}{\partial \rho^3}\right)_{S, \rho} = \rho_0$$

$$s = (\rho - \rho_0) / \rho_0.$$

The speed of sound becomes:

$$c = c_0^2 \left[ 1 + \left(\frac{B}{A}\right) s + \left(\frac{B}{2A}\right) s^2 + \dots \right]$$

and the equation of motion is, for the case where only the first two terms in the series are retained,

$$\xi = \frac{c_0^2}{(1 + \xi')^2 + B/A} \xi''.$$

Here it is seen that for the situation where the displacement amplitude cannot be neglected, the parameter B/A becomes a measure of the nonlinearity of the propagating medium.

A consequence of the propagation described by this relation in a fluid medium is that an originally sinusoidal, monochromatic wave becomes distorted as it propagates, harmonics are generated, and the amplitude of these harmonics is a function of the distance from the source(3). The harmonics have zero amplitude at the source, increase to a maximum value at a position from the source at which the effect of absorption processes balances such harmonic production and propagation occurs as under linear conditions well beyond this point.

The quantity B/A can be approximated as:

$$\frac{B}{A} = \frac{2 \rho_0 c_0^3}{\pi f} \left[ \frac{p_2(x)}{xp_1(o)} \right]^{-2}, \quad xp_1(o) \rightarrow 0$$

where  $p_1$  is the pressure amplitude of the fundamental,  $p_2$  is the pressure amplitude of the second harmonic, and  $x$  is propagation distance(1,4,6). A method which determines these quantities has the potential for yielding B/A for optically opaque media and for in vivo preparations.

Values of B/A have been obtained by measuring the harmonic content of pulses of sound at various distances from the source (3 MHz fundamental) and for varying concentrations in water of several biological macromolecules of interest(5). It has been found that B/A appears to increase nearly linearly with increasing concentration of proteins and to exhibit little dependence upon molecular weight (in the range  $10^2$  to  $10^6$  Daltons). These data suggest that the nonlinearity parameter increases with decreasing intersolute particle spacing. Thus it is clear that biological media exhibit significant nonlinear ultrasonic propagation features which must be studied in detail for deeper understanding leading to more sophisticated diagnostic and therapeutic procedures.

## REFERENCES

1. R.T. Beyer and S.V. Letcher, in "Physical Ultrasonics", Academic Press, New York (1969), pp. 202-224.
2. L.E. Kinsler and A.R. Frey, in "Fundamentals of Acoustics", John Wiley, New York (1962), 2nd ed., pp. 108-113.
3. W.J. Fry and F. Dunn, Ultrasound: analysis and experimental methods in biological research, in "Physical Techniques in Biological Research", W.L. Nastuk, ed., Academic Press, New York (1962), Vol. 4, Ch. 6, pp. 261-394.
4. E. Fubni-Ghiron, Anomalie nella propagazione di onde acustiche di grande ampiezza, Alta Frequenza 4:530-581 (1935).
5. W.K. Law, L.A. Frizzell and F. Dunn, Ultrasonic determination of the nonlinearity parameter B/A for biological media, J. Acoust. Soc. Am. 69:4 (1981).
6. L. Adler and E.A. Hiedmann, Determination of the nonlinearity parameter B/A for water and m-xylene, J. Acoust. Soc. Am. 34:410-411 (1962).

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## PROPAGATION OF ULTRASOUND IN SOLUTIONS OF BIOLOGICAL SUBSTANCES

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The values of ultrasound velocity and absorption coefficient in solutions are defined by different molecular interactions. There is a large amount of literature on the investigation of molecular characteristics of solutions of biological substances by ultrasonic measurements. The dominant part of this paper is devoted to the measurements of the frequency dependences of ultrasound absorption and the investigation of fast relaxation processes. Substantial works are related to the study of solutions by ultrasound velocity measurements. But such a proportion is a result of the absence of adequate velocity measurement methods and not due to the fact that the absorption coefficient is more informative than velocity of ultrasound about the characteristics of a solution.

The purpose of this short review is to show the relations between acoustic characteristics of biological solutions and their molecular properties, with emphasis on the ultrasound velocity. There are two reasons for such intentions:

- (a) The possibilities of ultrasound velocity measurements in investigation of biological substances in solutions are much less known, and
- (b) The value of ultrasound velocity reflects to a greater extent the molecular characteristics of solutions of biological substances than absorption.